

Radio and mm astronomy

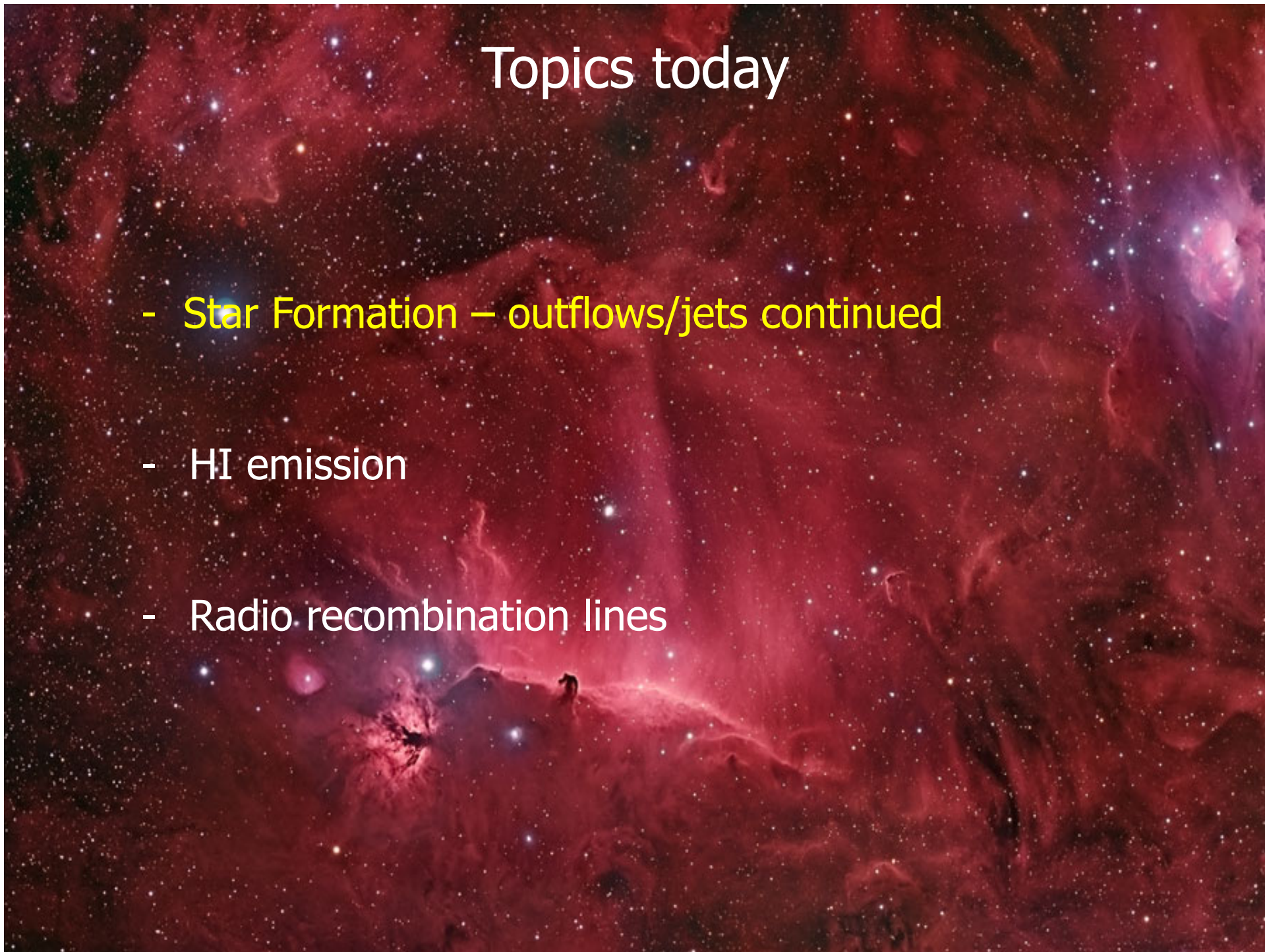
Wintersemester 2012/2013
Henrik Beuther & Hendrik Linz

<i>16.10 Introduction & Overview</i>	<i>(HL & HB)</i>
<i>23.10 Emission mechanisms, physics of radiation</i>	<i>(HB)</i>
<i>30.10 Telescopes – single-dishes</i>	<i>(HL)</i>
<i>06.11 Telescopes – interferometers</i>	<i>(HB)</i>
<i>13.11 Instruments – continuum radiation</i>	<i>(HL)</i>
<i>20.11 Instruments – line radiation</i>	<i>(HB)</i>
<i>27.11 Continuous radiation (free-free, synchrotron, dust, CMB)</i>	<i>(HL)</i>
<i>04.12 Line radiation</i>	<i>(HB)</i>
<i>11.12 Radiation transfer</i>	<i>(HL)</i>
<i>18.12 Effelsberg Excursion</i>	
<i>Christmas break</i>	
<i>08.01 Molecules and chemistry</i>	<i>(HL)</i>
<i>15.01 Physics and kinematics</i>	<i>(HB)</i>
<i>22.01 Applications</i>	<i>(HL)</i>
29.01 Applications	(HB)
<i>05.02 last week, no lecture</i>	

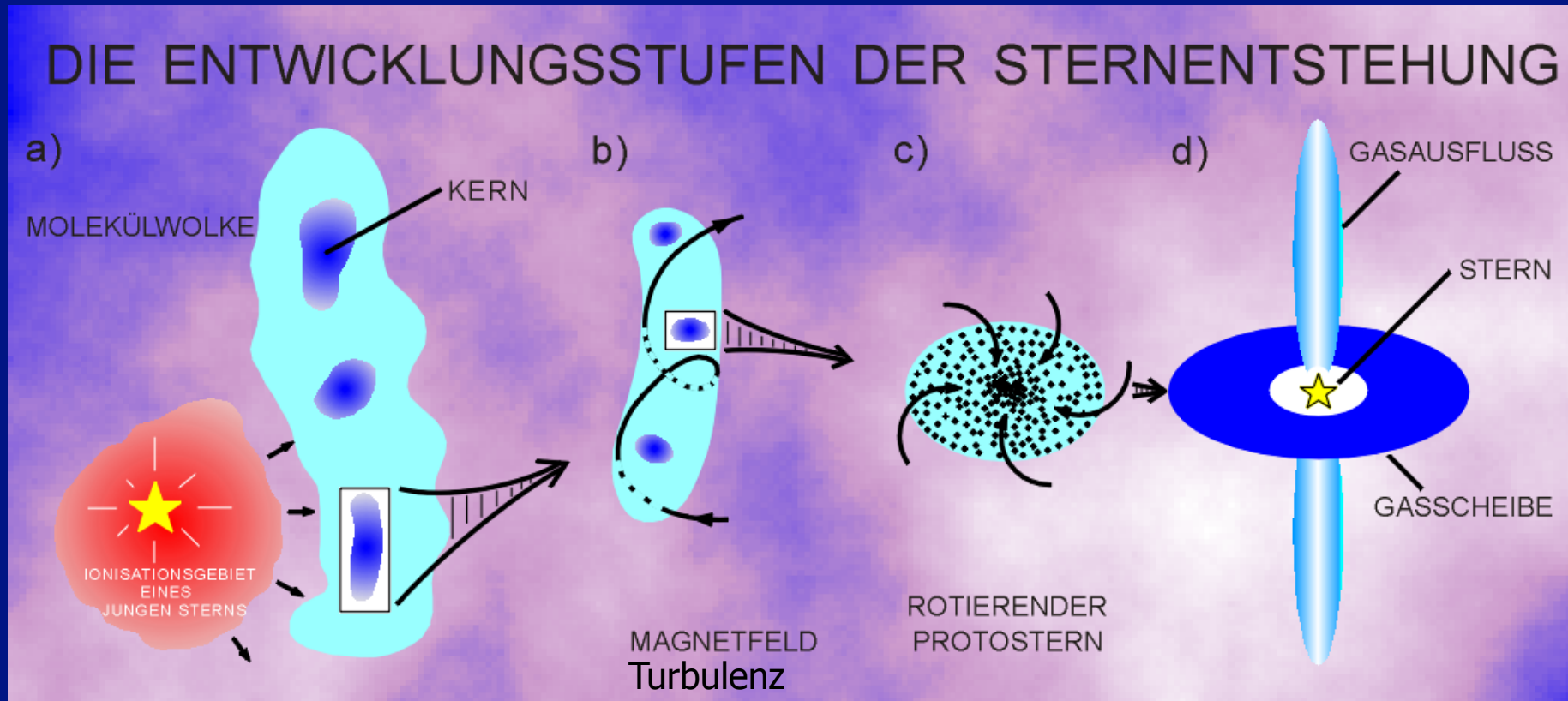
More Information and the current lecture files: http://www.mpia.de/homes/beuther/lecture_ws1213.html
beuther@mpia.de, linz@mpia.de

Topics today

- Star Formation – outflows/jets continued
- HI emission
- Radio recombination lines



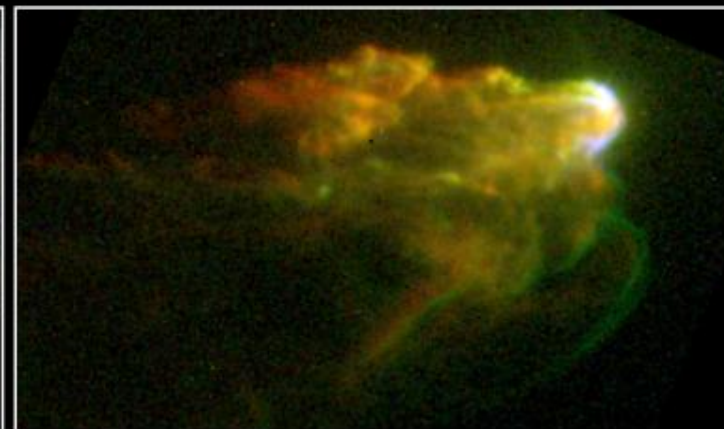
Sternentstehungsparadigma



Time scales: Main accretion phase $\sim 500\,000$ years
Pre-main-sequence evolution $\sim 2e6$ years

Discovery of outflows I

Herbig 1950, 1951; Haro 1952, 1953



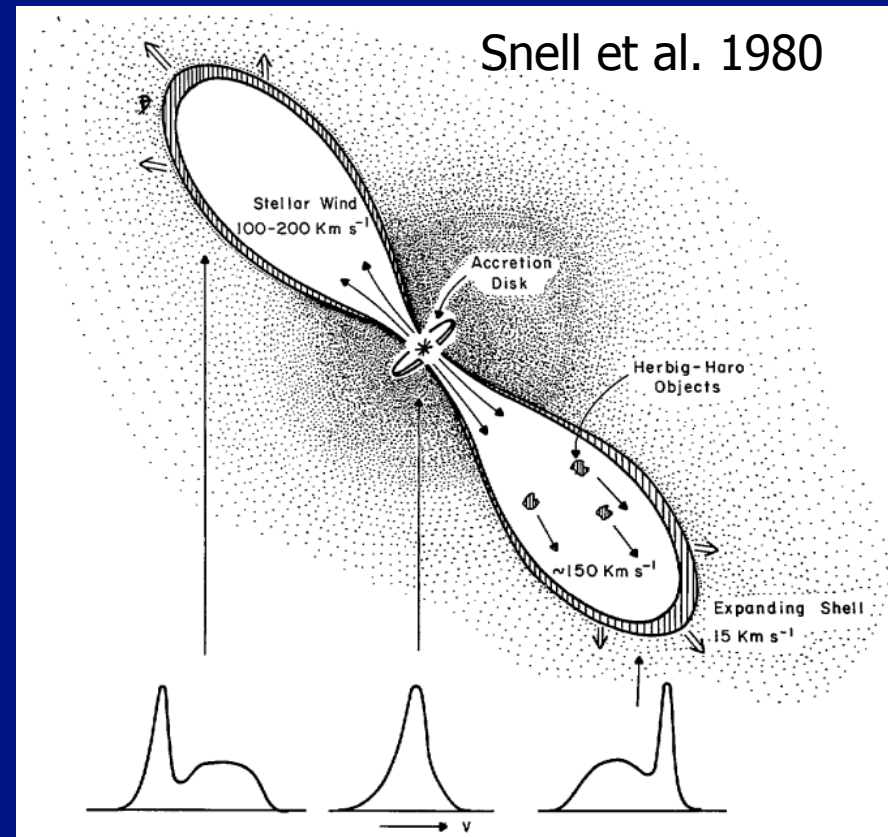
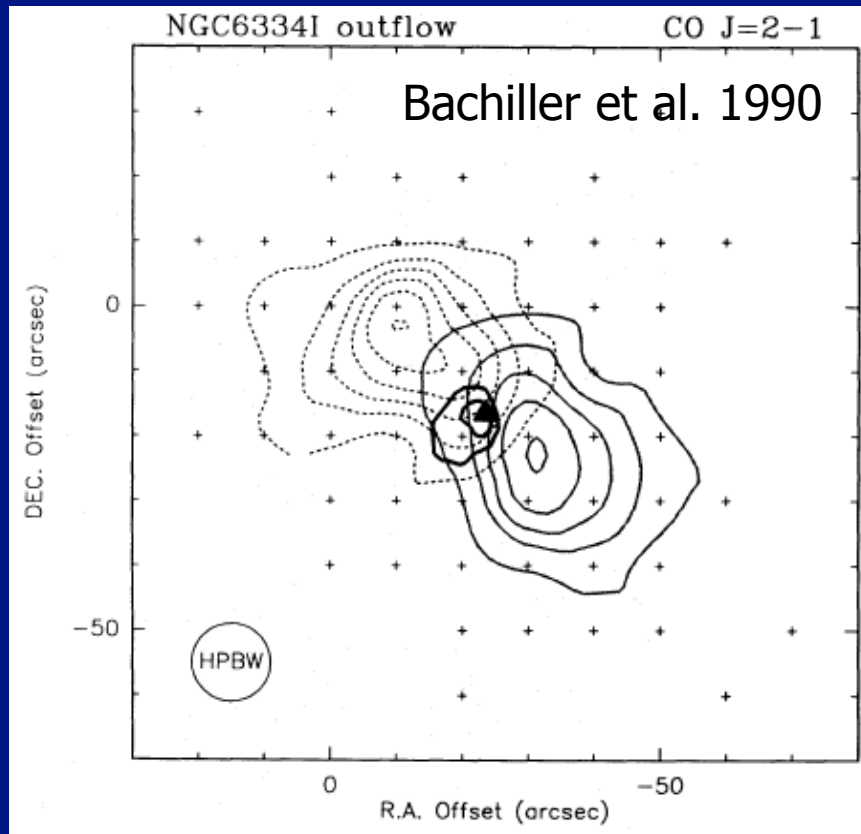
Jets from Young Stars · HH1/HH2

HST · WFPC2

PRC95-24c · ST ScI OPO · June 6, 1995 · J. Hester (AZ State U.), NASA

Initially thought to be embedded protostars but soon spectra were recognized as caused by shock waves → jets and outflows indicated

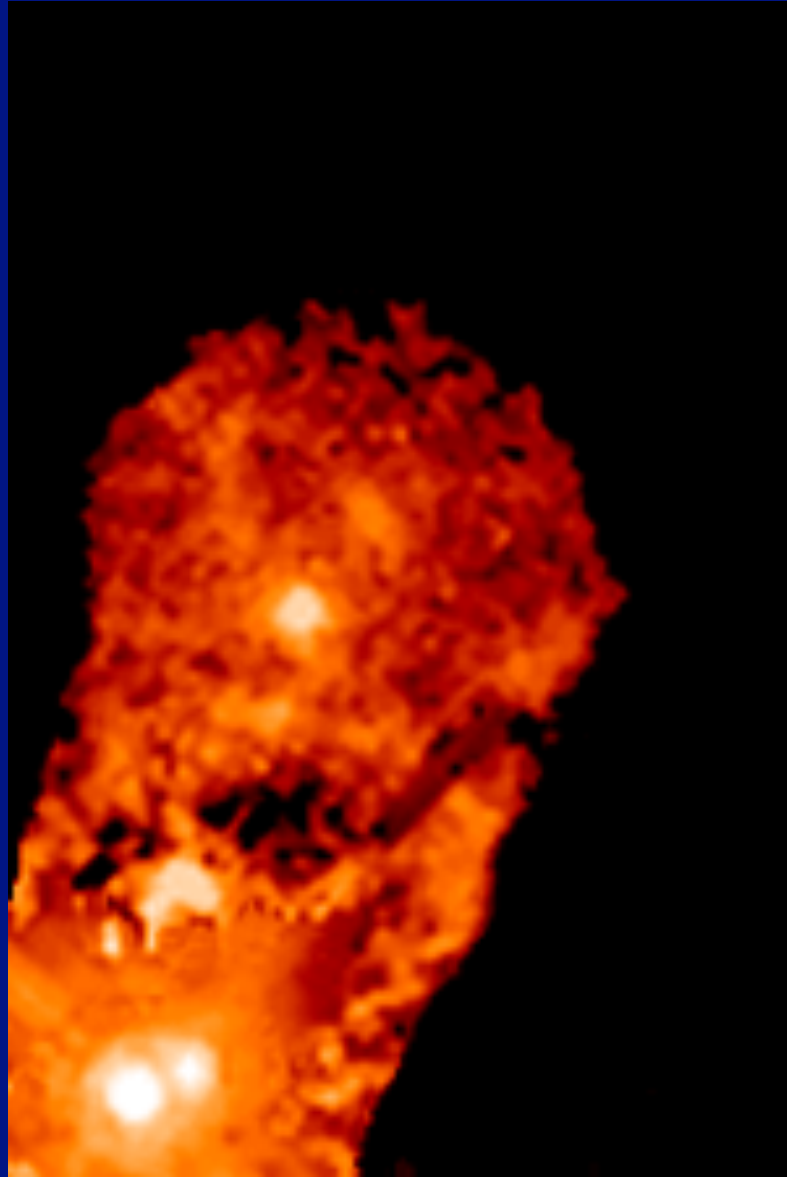
Discovery of outflows II



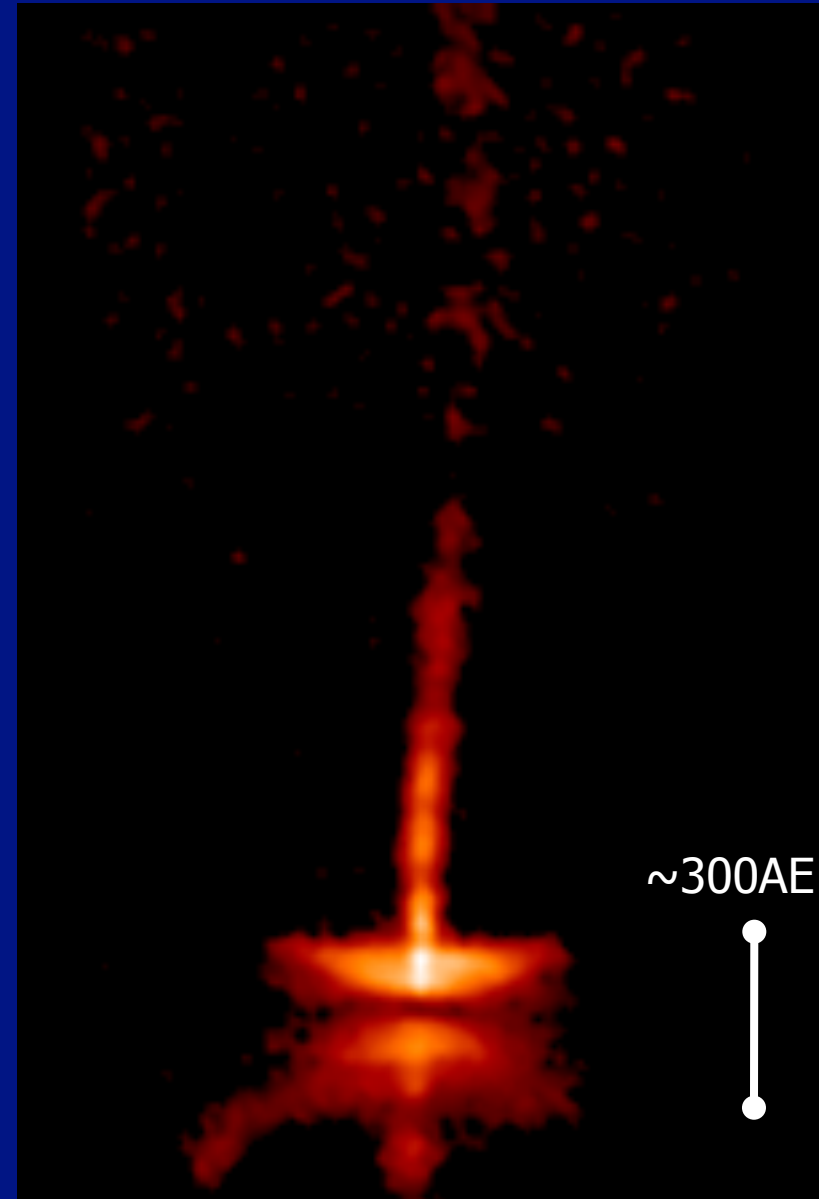
- In the mid to late 70th, first CO non-Gaussian line wing emission detected (Kwan & Scoville 1976).
- Bipolar structures, extremely energetic, often associated with HH objects

Eigenbewegungen von Jets

XZ Tauri



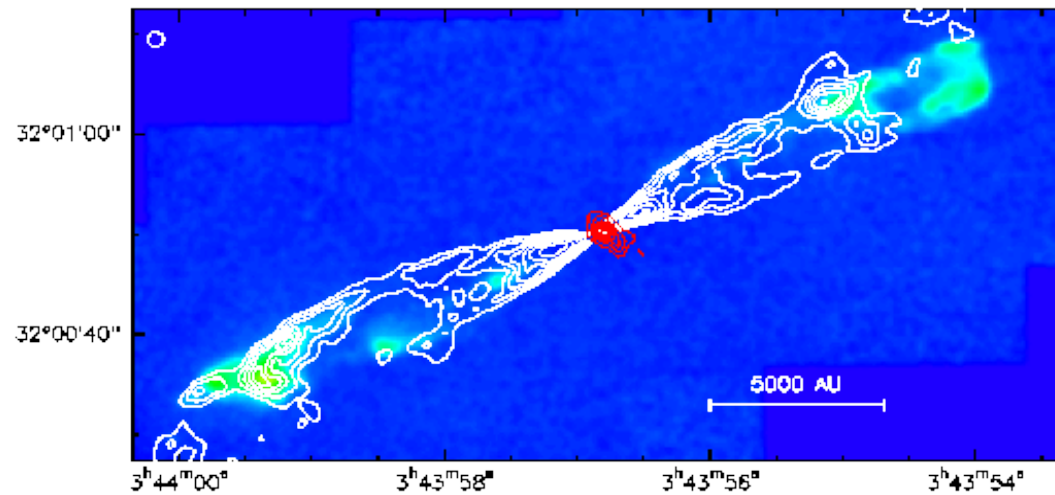
HH30



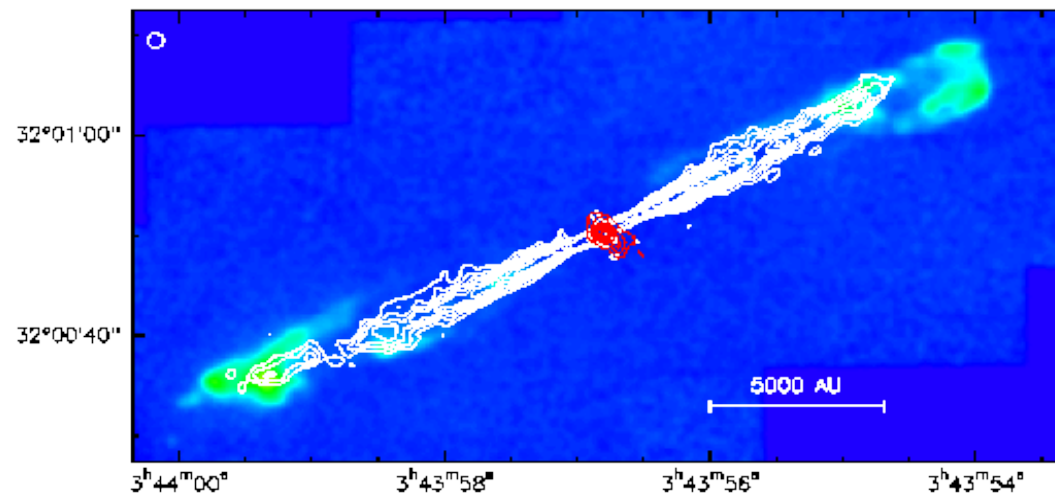
The prototypical molecular outflow HH211

HH211, Gueth et al. 1999

H₂ 2.12 μm (colors) + CO J=2-1 v<10 km/s (white) + continuum 1.3 mm (red)



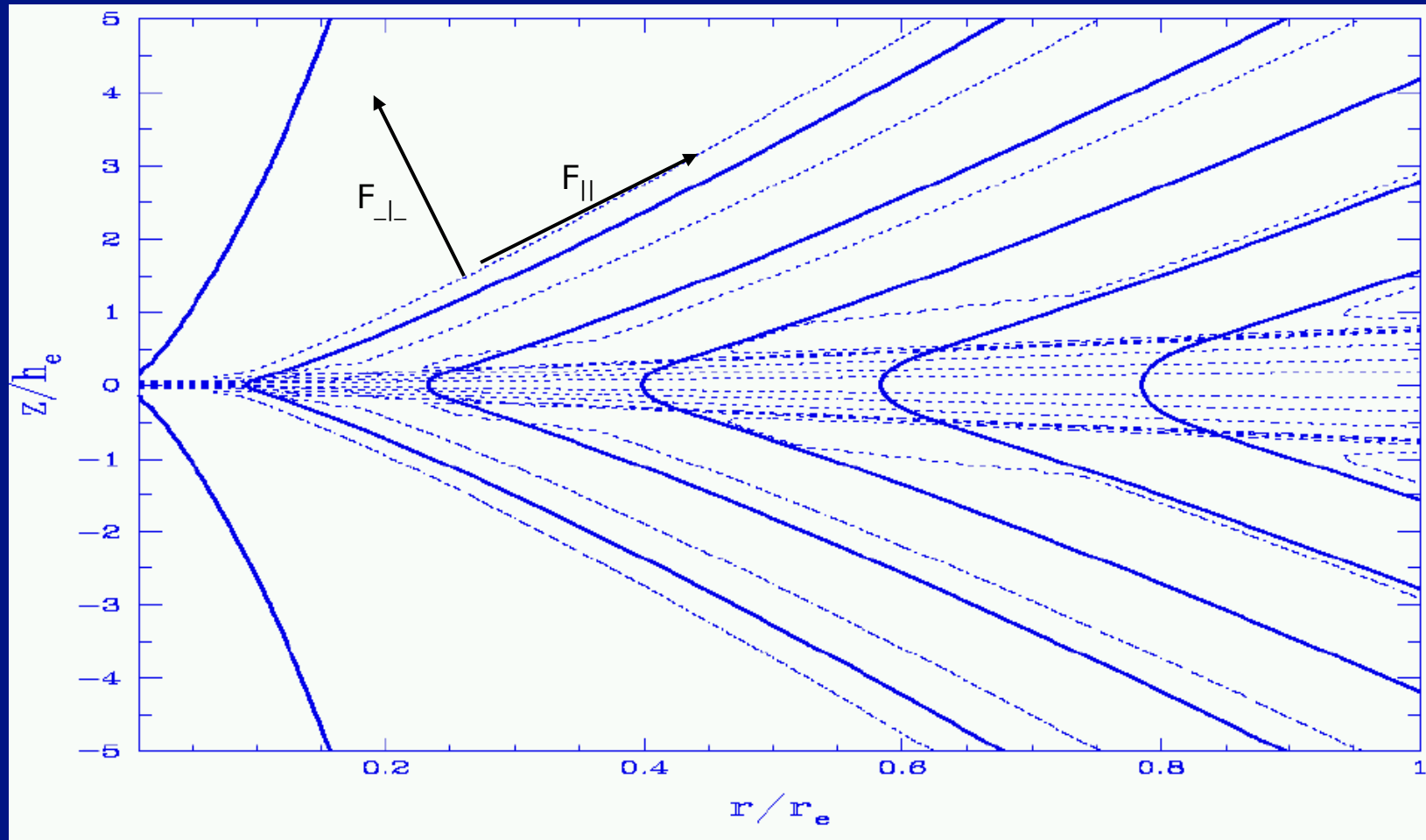
H₂ 2.12 μm (colors) + CO J=2-1 v>10 km/s (white) + continuum 1.3 mm (red)



Jet launching from accretion disks

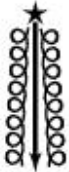






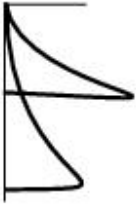



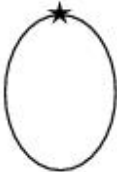

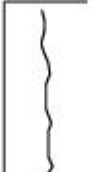






“magnetic accretion-ejection structures” (Ferreira et al 1995-1997):

- 1) disk material diffuses across magnetic field lines,
- 2) is lifted upwards by MHD forces, then
- 3) couples to the field and 4) becomes accelerated magnetocentrifugally and 5) collimated



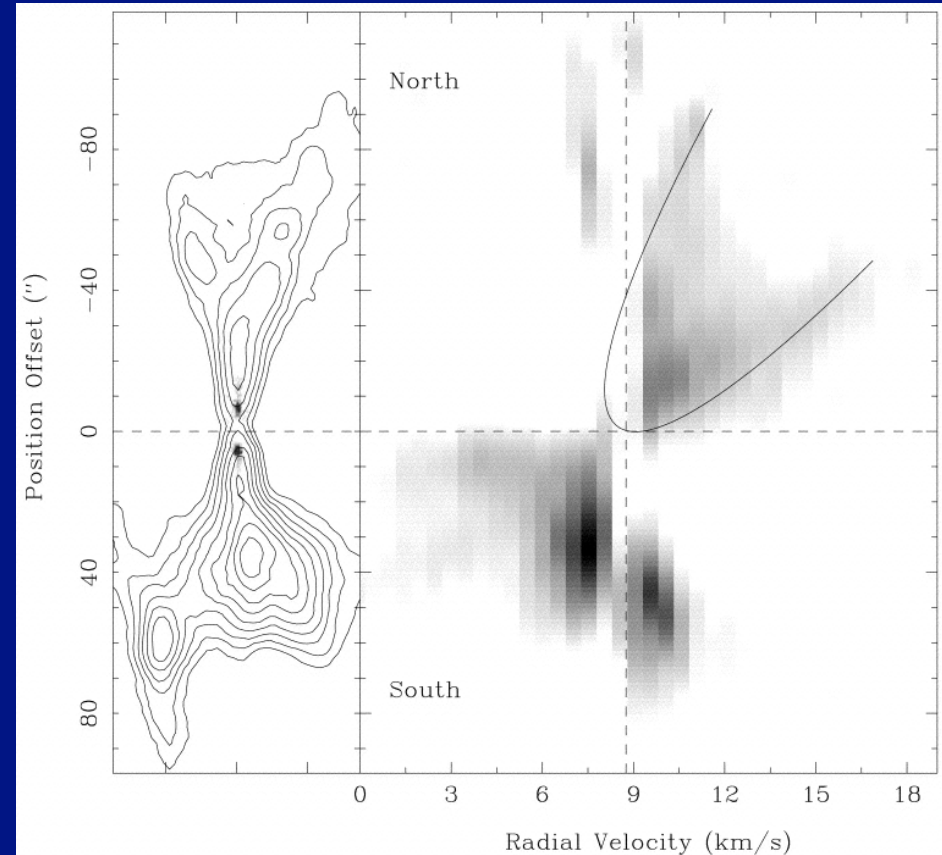
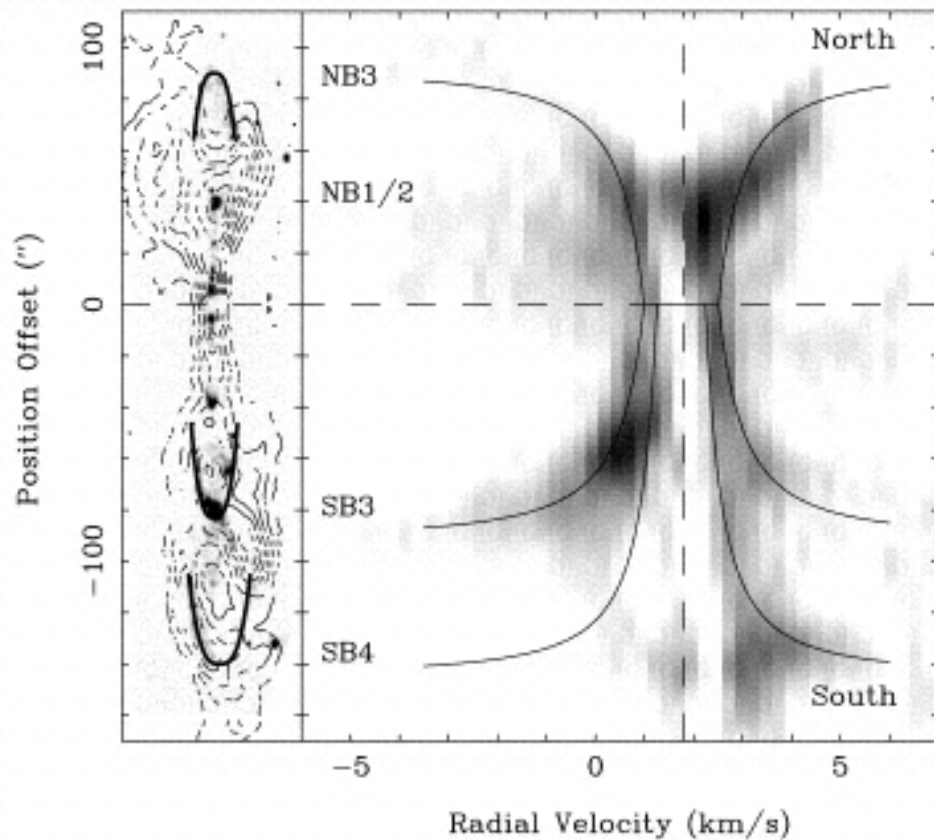
Magnetic field lines (thick)
and streamlines (dashed)

Outflow entrainment models

Molecular outflow properties predicted by different models					
Model	Wind	Predicted property of molecular outflow along axis			
		Morphology	Velocity	Temperature	Momentum ^a
Turbulent Jet					
Jet Bow Shock					
Wide-angle Wind					
Circulation					

^a Assuming an underlying density distribution of r^{-1} to r^{-2} .

Collimation and pv-structure



HH212: consistent with jet-driving

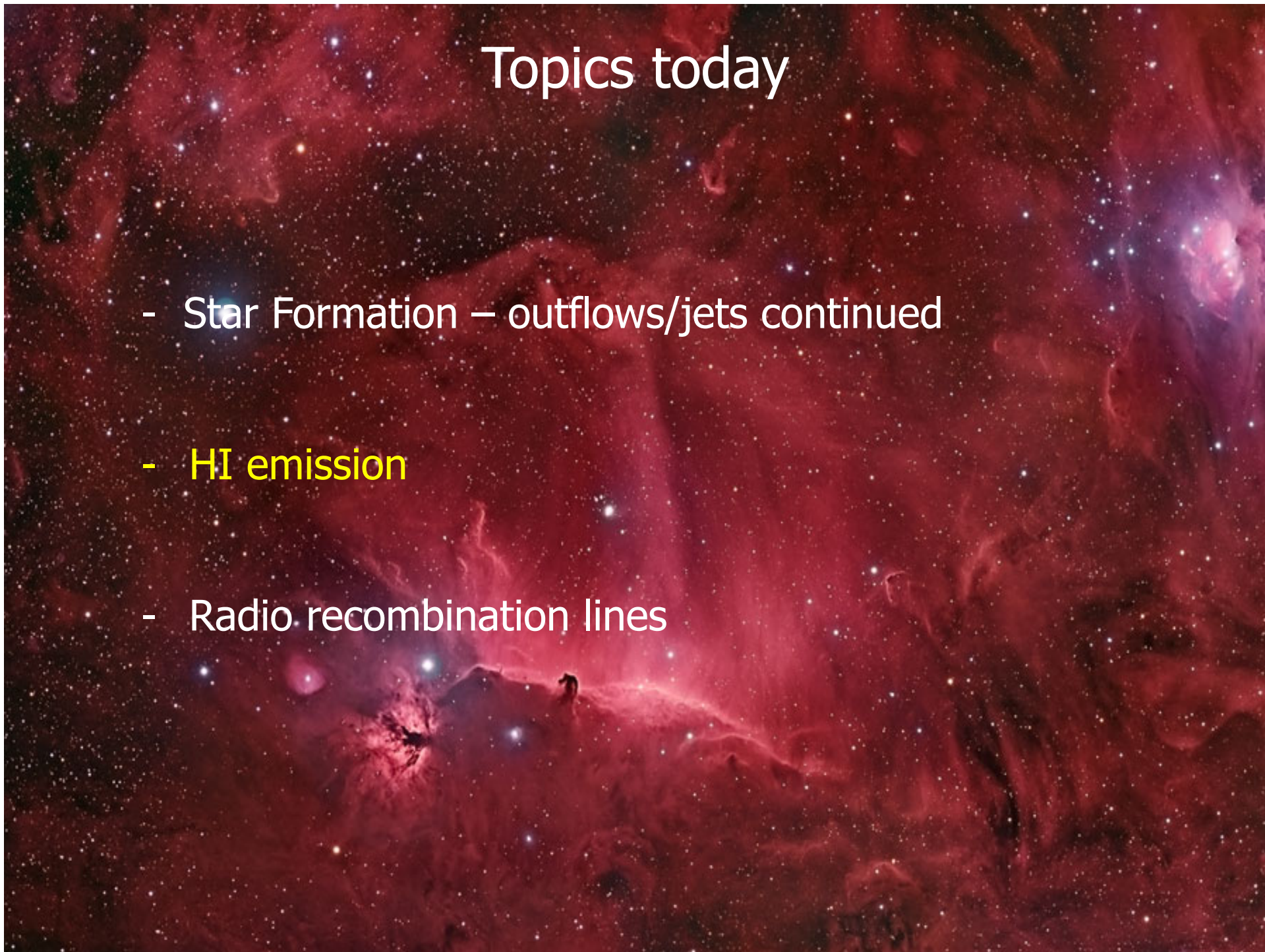
VLA0548: consistent with wind-driving

- pv-structure of jet- and wind-driven models very different
- Often Hubble-law observed --> increasing velocity with increasing distance from the protostar

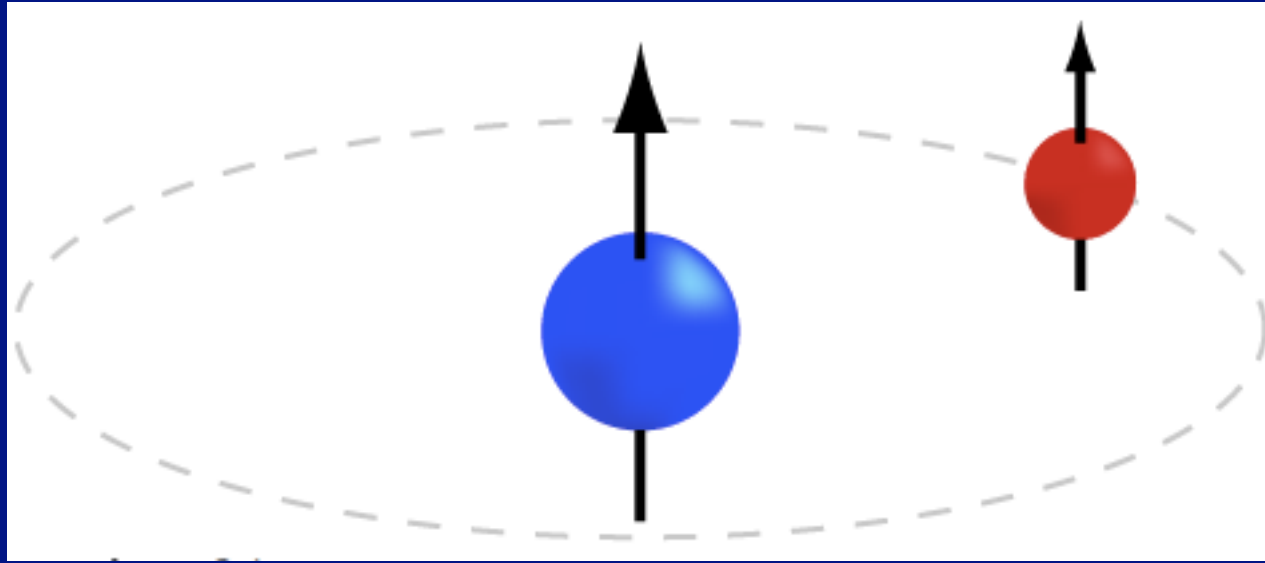
Lee et al. 2001

Topics today

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- **HI emission**
- Radio recombination lines



Atomic hydrogen HI



Two energy levels from the magnetic interaction between proton and e^- spin.
→ Spin flip causes line emission at 1.4GHz or 21cm

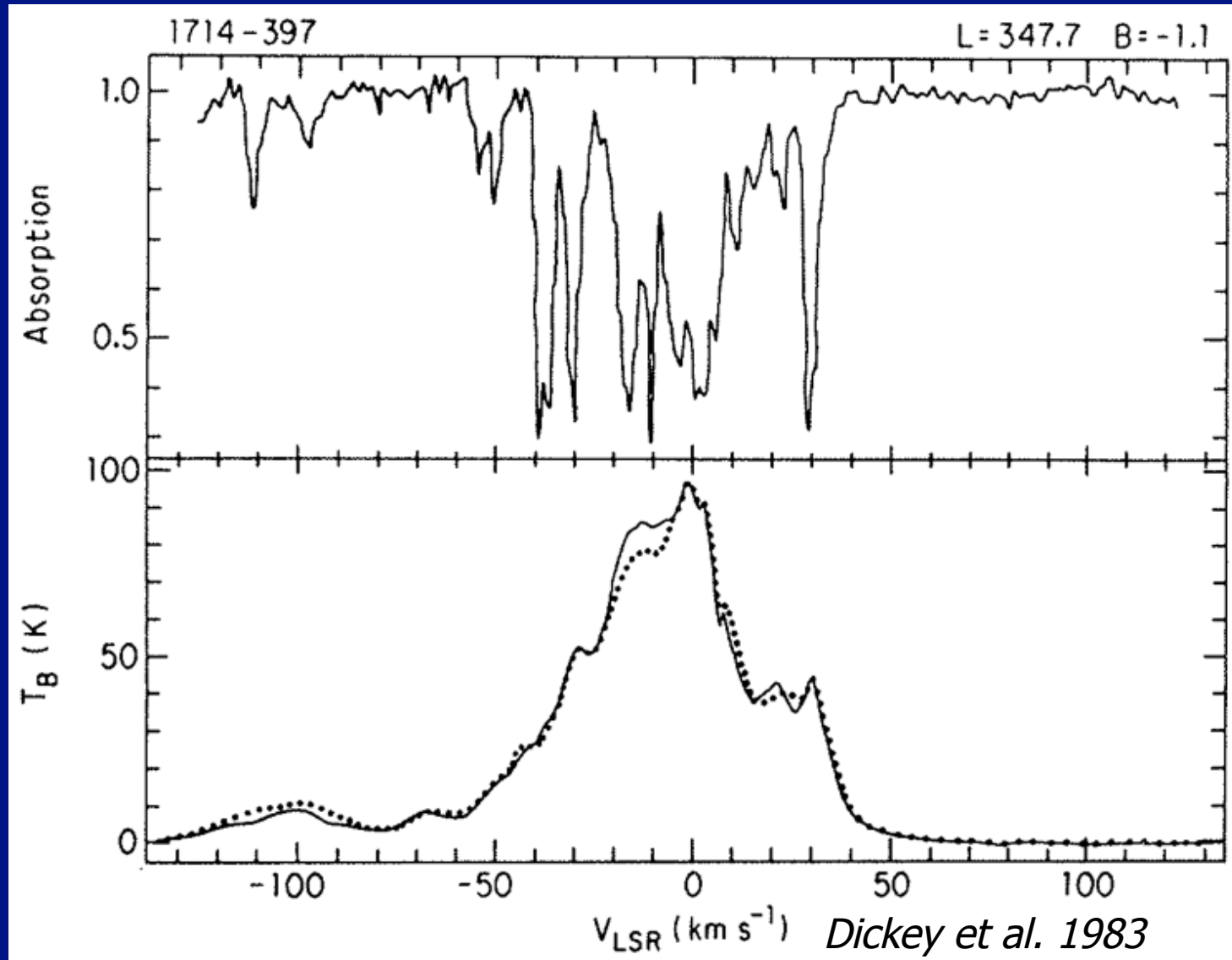
Einstein A coefficient $A_{10} \approx 2.85 \times 10^{-15} \text{ s}^{-1}$ very low!

That results in a half-time: $\tau_{1/2} = A_{10}^{-1} \approx 3.5 \times 10^{14} \text{ s} \approx 11 \text{ million years}$

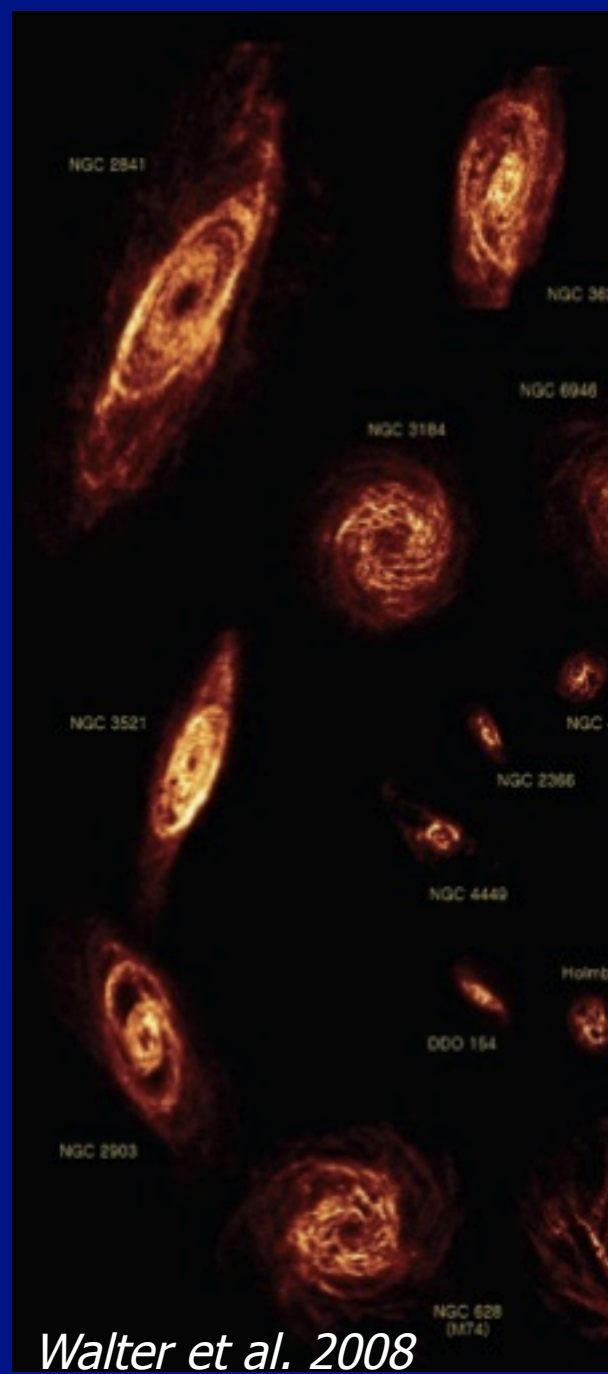
This implies critical densities $\ll 1 \text{ cm}^{-3}$ and collisional equilibrium can easily be maintained.

→ Early detections were only possible at low densities of ISM!

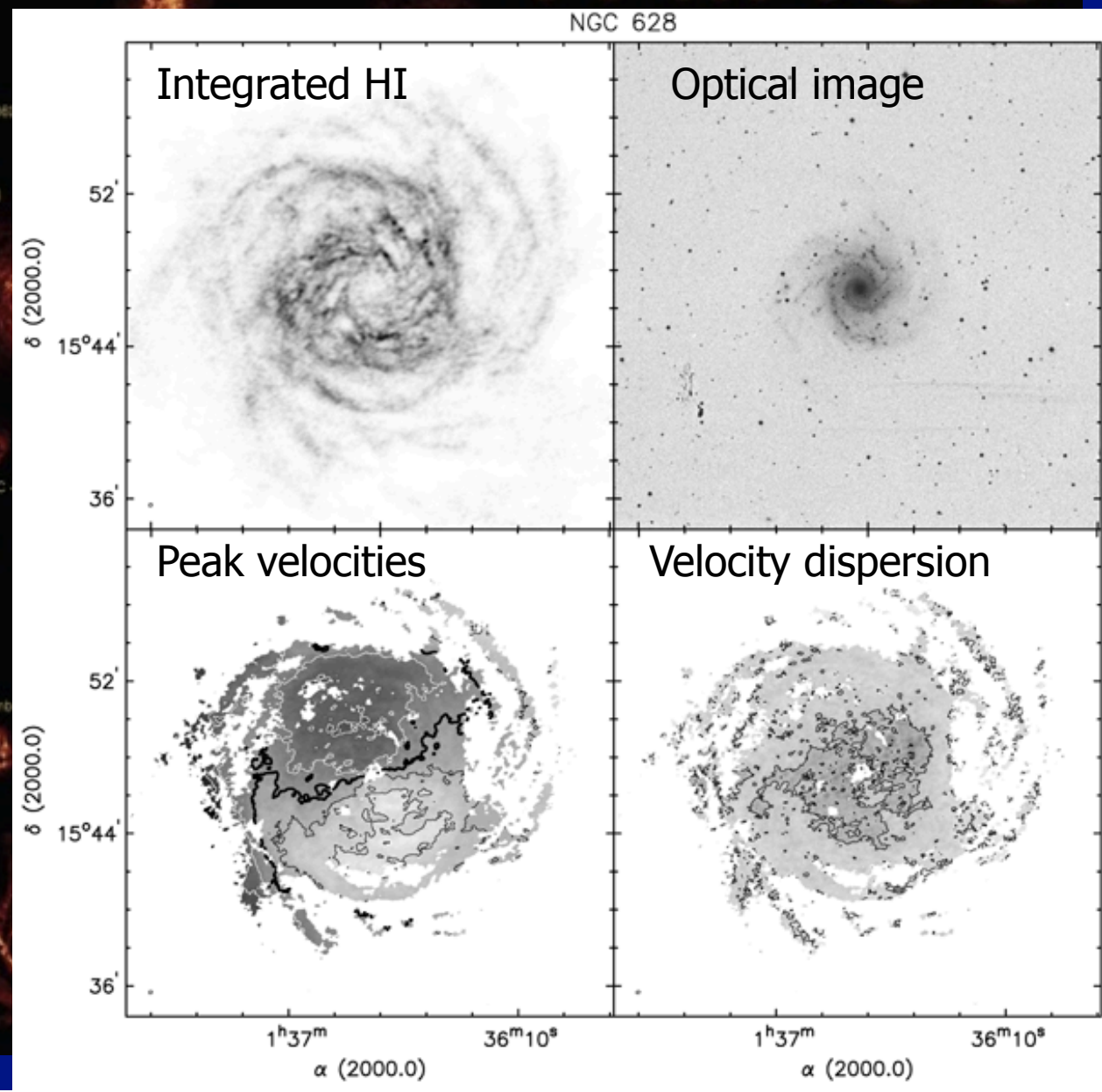
HI in emission and absorption



THINGS

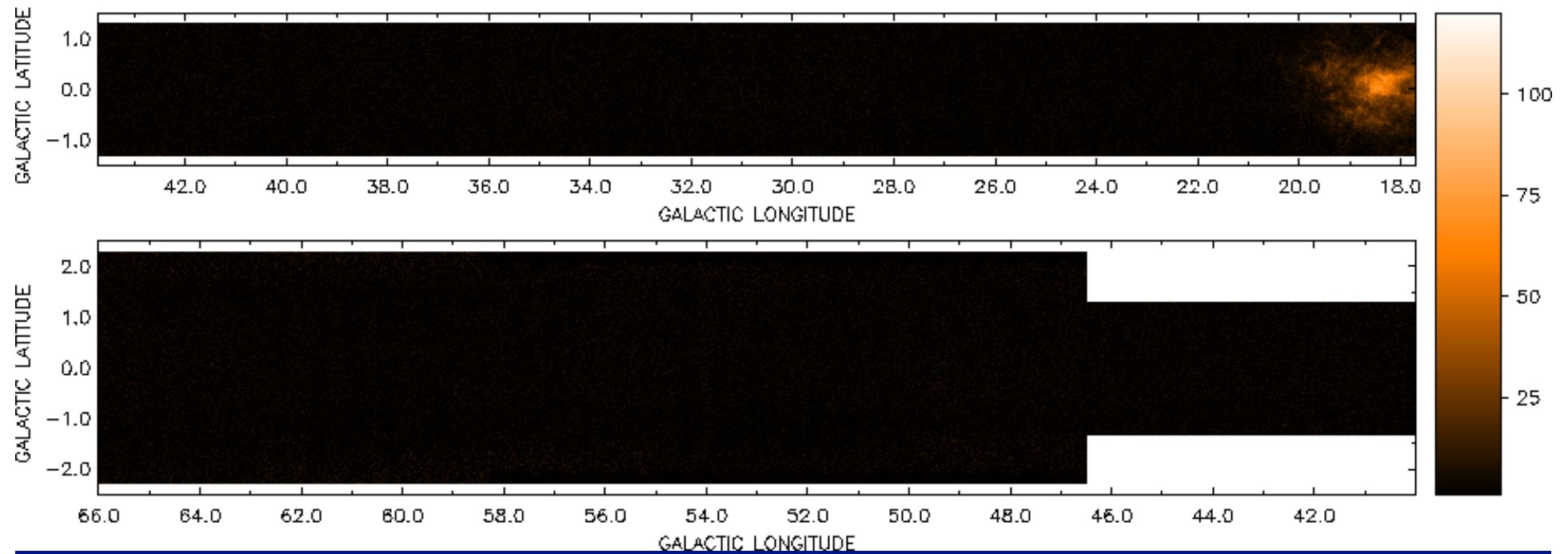


Walter et al. 2008



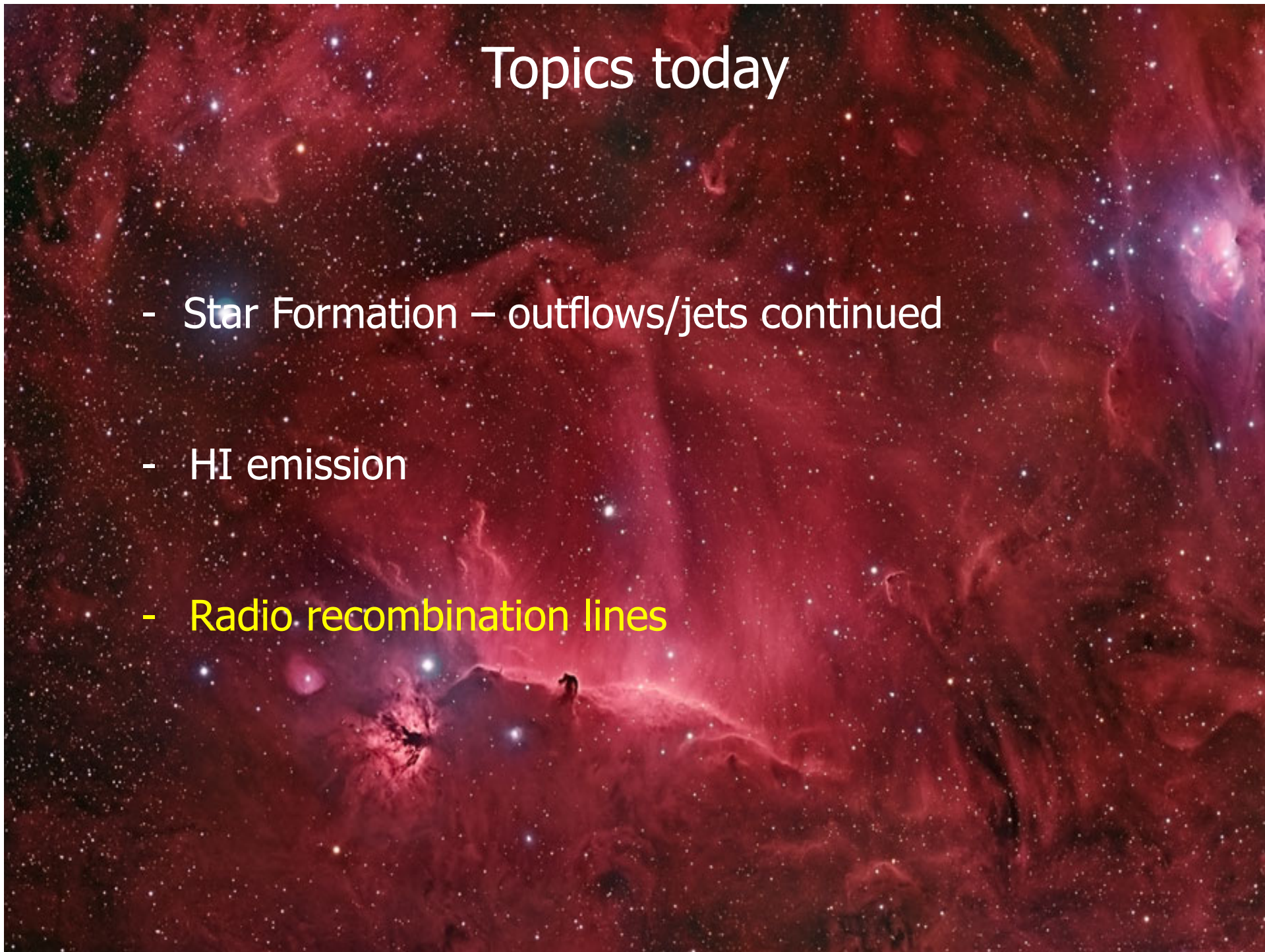
HI in the Milky Way

GALACTOCENTRIC RADIUS 3.00 kpc

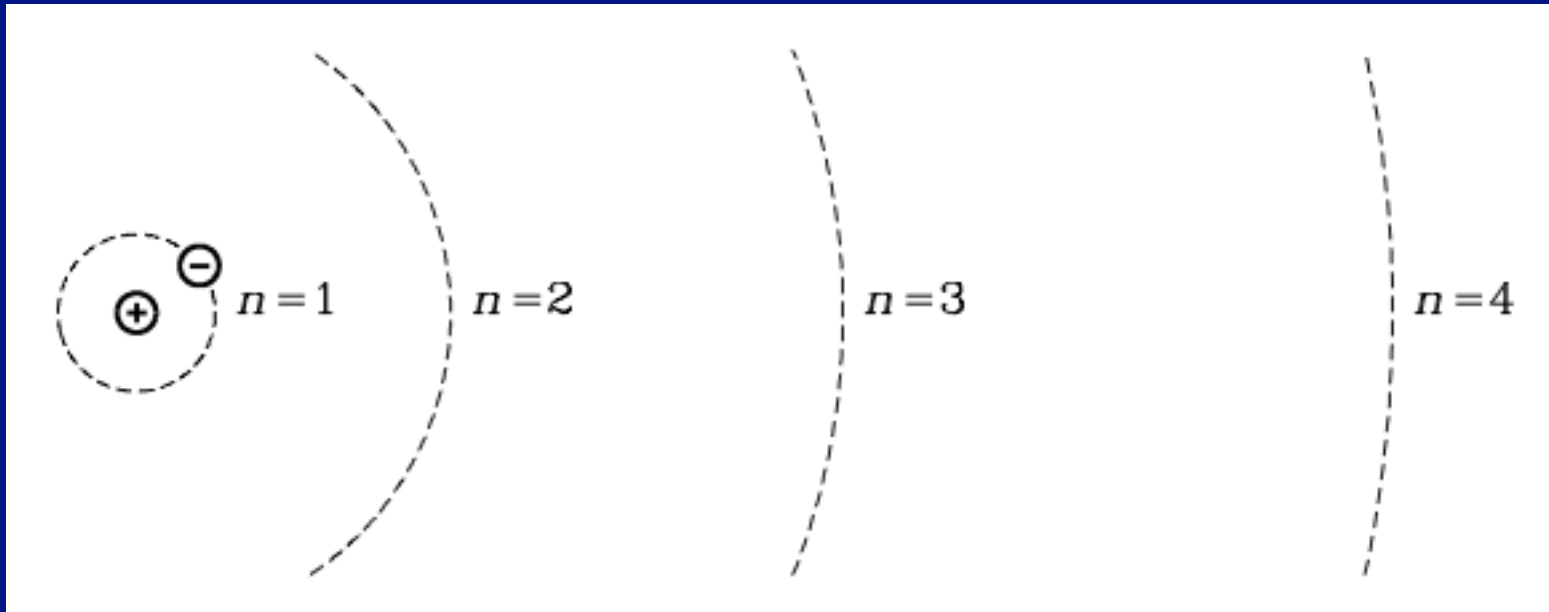


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Recombination lines I



- After recombination, the electron can cascade down the “shells”
→ these transitions are called recombination lines

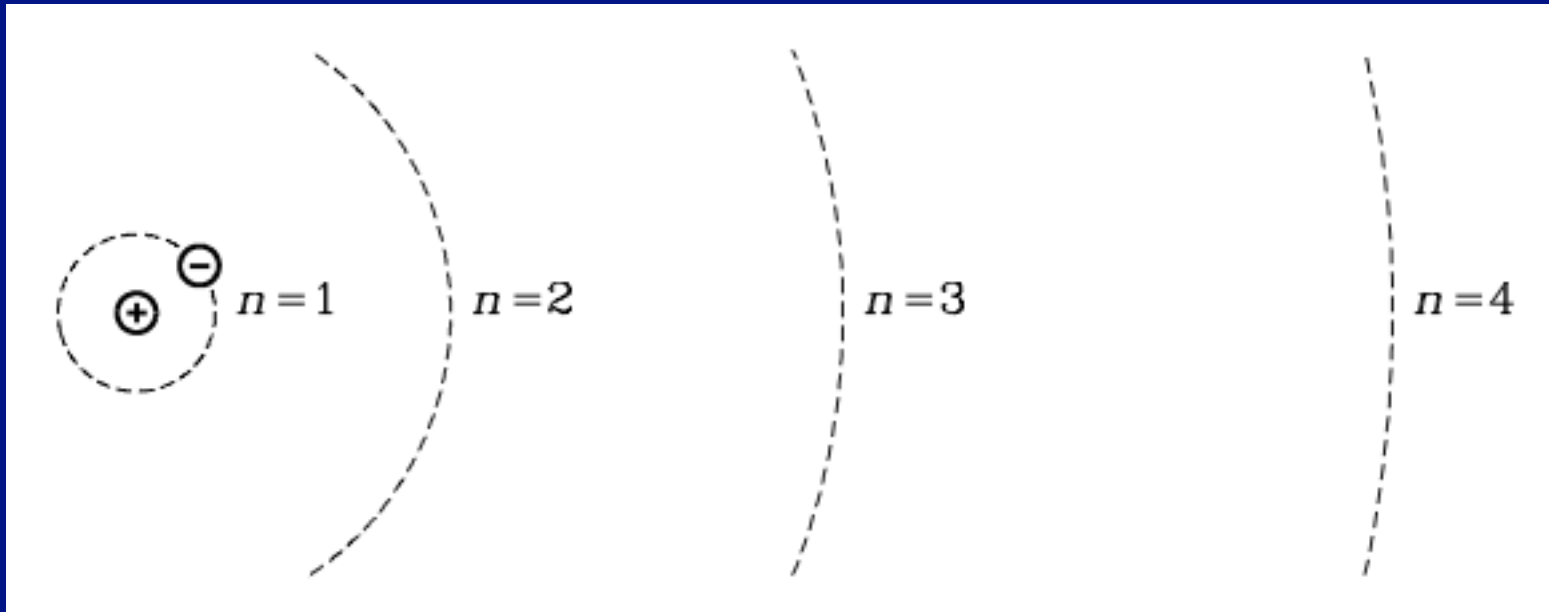
- De-Broglie wavelength $\lambda = \frac{h}{p} = \frac{h}{m_e v}$ with m_e = mass electron

- Only orbits with circumferences of an integer number of n correspond to standing waves and are permitted:

$$2\pi a_n = n\lambda = \frac{nh}{m_e v}$$

with a_n = Bohr radius

Recombination lines II



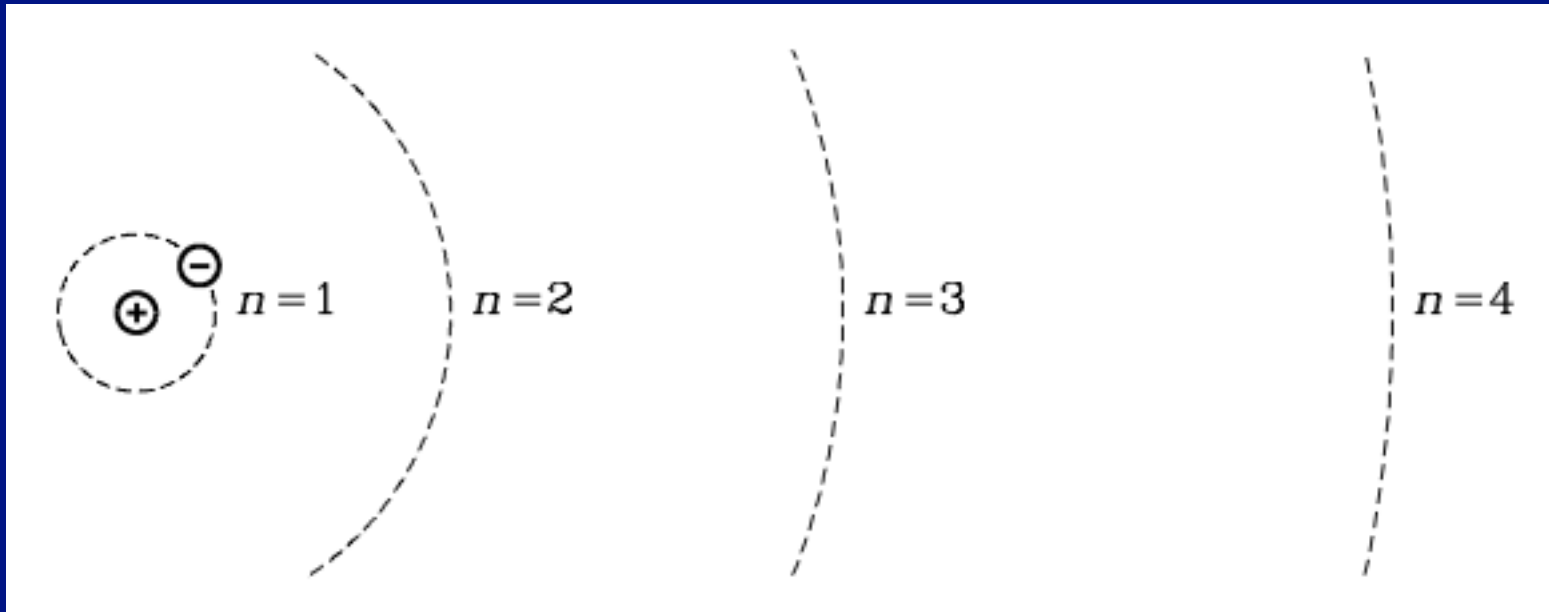
$$\rightarrow a_n = \frac{nh}{2\pi m_e v} = \frac{n\hbar}{m_e v}$$

- Relation between a and v can be derived from equilibrium between coulomb force and centrifugal force:

$$\frac{e^2}{a_n^2} = \frac{m_e v^2}{a_n}$$

$$\rightarrow v^2 = \frac{e^2}{m_e a_n} \quad \text{or} \quad a_n = \frac{n^2 \hbar^2}{m_e e^2}$$

Recombination lines III



- Bohr radius of hydrogen atom:

→ radius depends on n^2 !

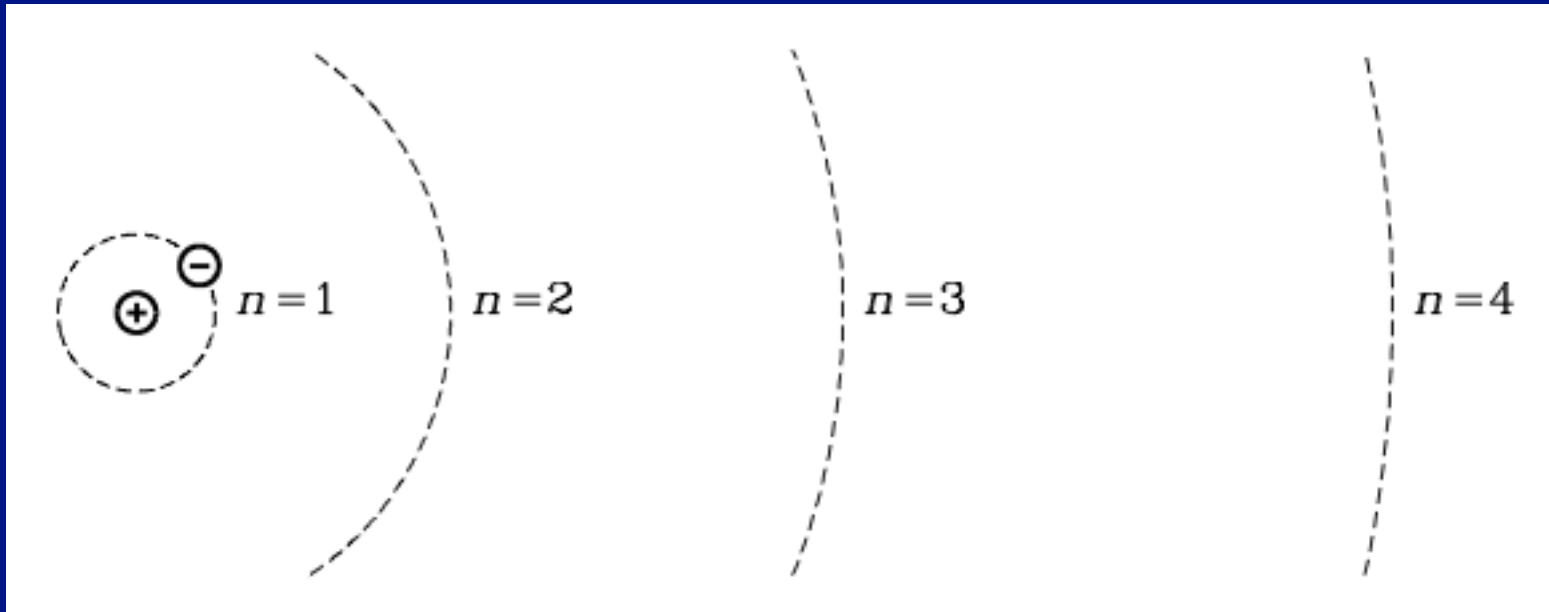
$$a_n = \frac{\hbar^2}{m_e e^2} n^2 \approx 0.53 \times 10^{-8} \text{ cm} \times n^2$$

- For $n=1 \rightarrow a_1 \sim 0.53 \times 10^{-8} \text{ cm}$

- For $n=100 \rightarrow a_{100} \sim 10^{-4} \text{ cm} = 1 \mu\text{m}$

→ nearly macroscopic sizes, larger than for example many viruses

Recombination lines IV



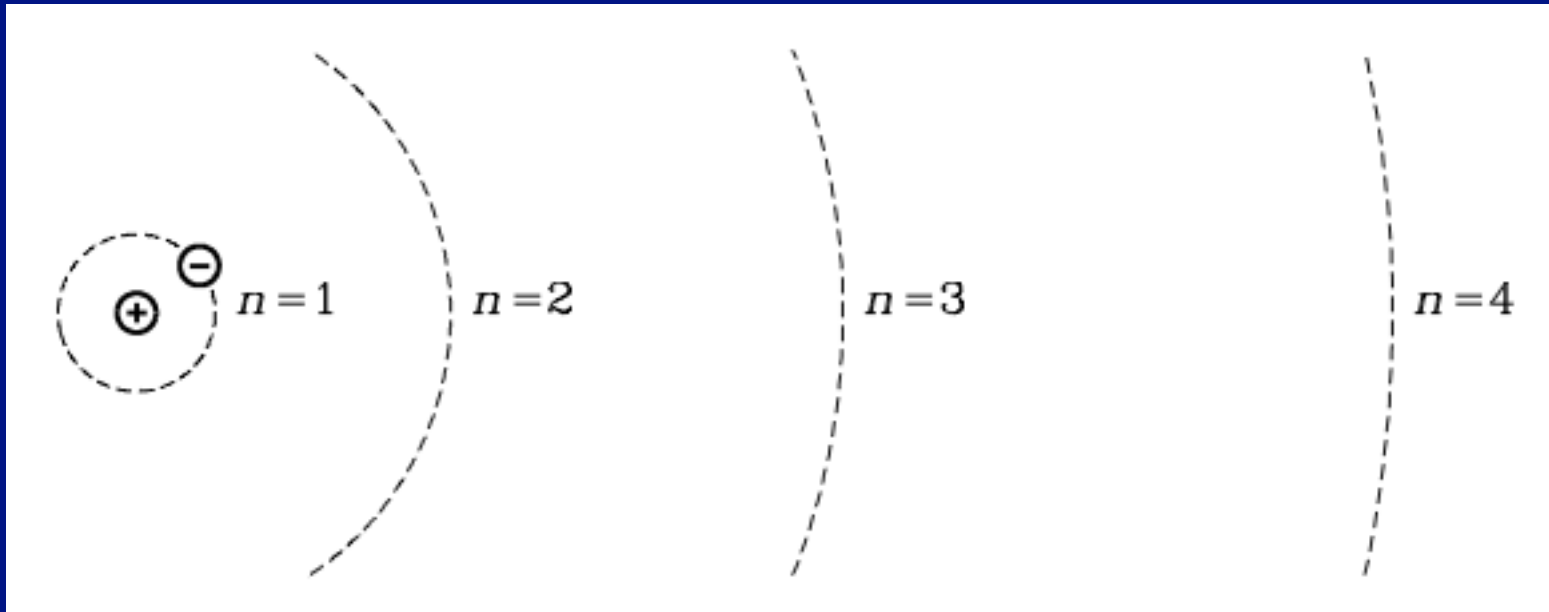
- Energy in n^{th} level sum of kinetic and potential/coulomb energy

$$E_n = T + V = -T = V/2 = -\frac{e^2}{2a_n} = -\frac{e^2 m_e e^2}{2n^2 \hbar^2} = -\frac{m_e e^4}{2\hbar^2 n^2}$$

→ energy change ΔE between levels n and $n+\Delta n$

$$\Delta E = \frac{m_e e^4}{2\hbar^2} \left[\frac{1}{n^2} - \frac{1}{(n + \Delta n)^2} \right] = h\nu$$

Recombination lines V



→ frequency:
$$\nu = \left(\frac{2\pi^2 m_e e^4}{h^3 c} \right) c \left[\frac{1}{n^2} - \frac{1}{(n + \Delta n)^2} \right]$$

with Rydberg constant:
$$R_\infty \equiv \left(\frac{2\pi^2 m_e e^4}{h^3 c} \right) = 1.09737312 \dots \times 10^5 \text{ cm}^{-1}$$

(infinite sign because mass of nucleus is considered large compared to electron mass in this analysis)

- Labeling for lines use lower level n and Δn

$$\alpha \text{ for } \Delta n = 1, \beta \text{ for } \Delta n = 2, \gamma \text{ for } \Delta n = 3$$

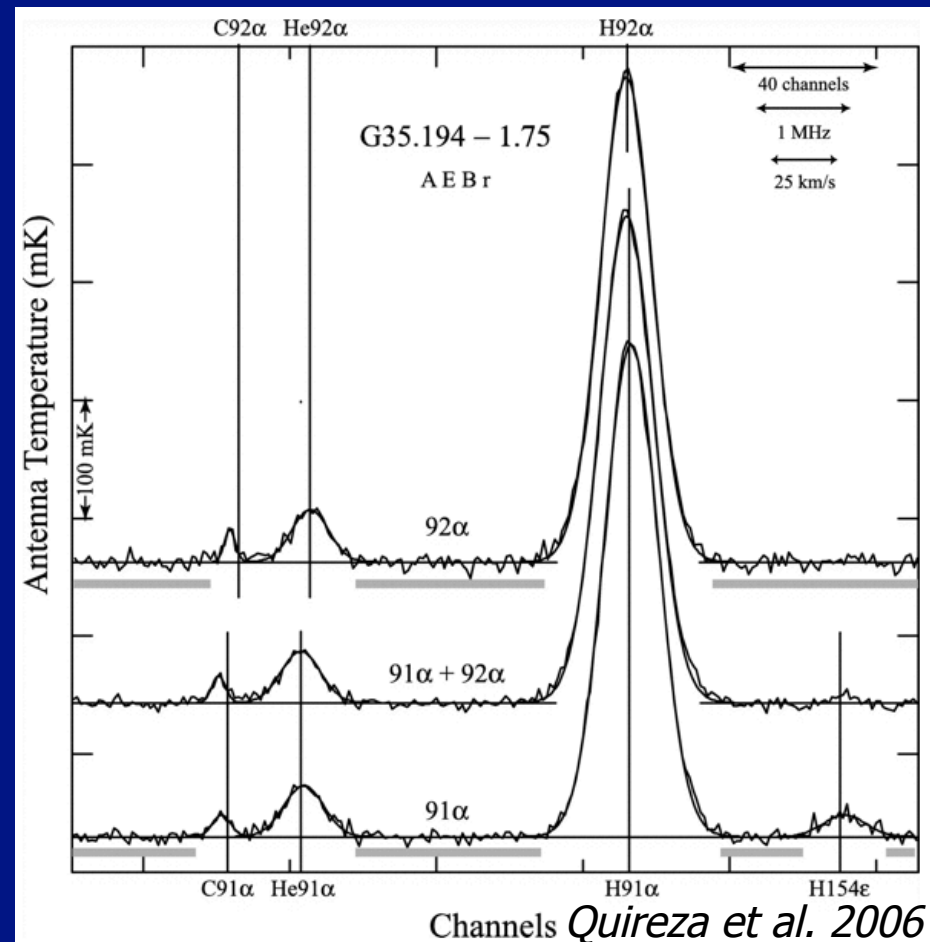
Recombination lines VI

- Low-n lines in the UV, optical and infrared, Lyman-alpha at 121.6nm
- Other famous series: n=2 Balmer, n=3 Paschen, n=4 Brackett
n=5 Pfund, n=6 Humphreys (near- to mid-infrared)
- In the radio: H109 α , 1965 by P. Mezger

- Analysis with finite nuclear mass:
Rydberg constant:

$$R_M \equiv R_\infty \left(1 + \frac{m_e}{M} \right)^{-1}$$

→ This implies that the frequencies of the heavier atoms' recombination lines are very close to hydrogen.

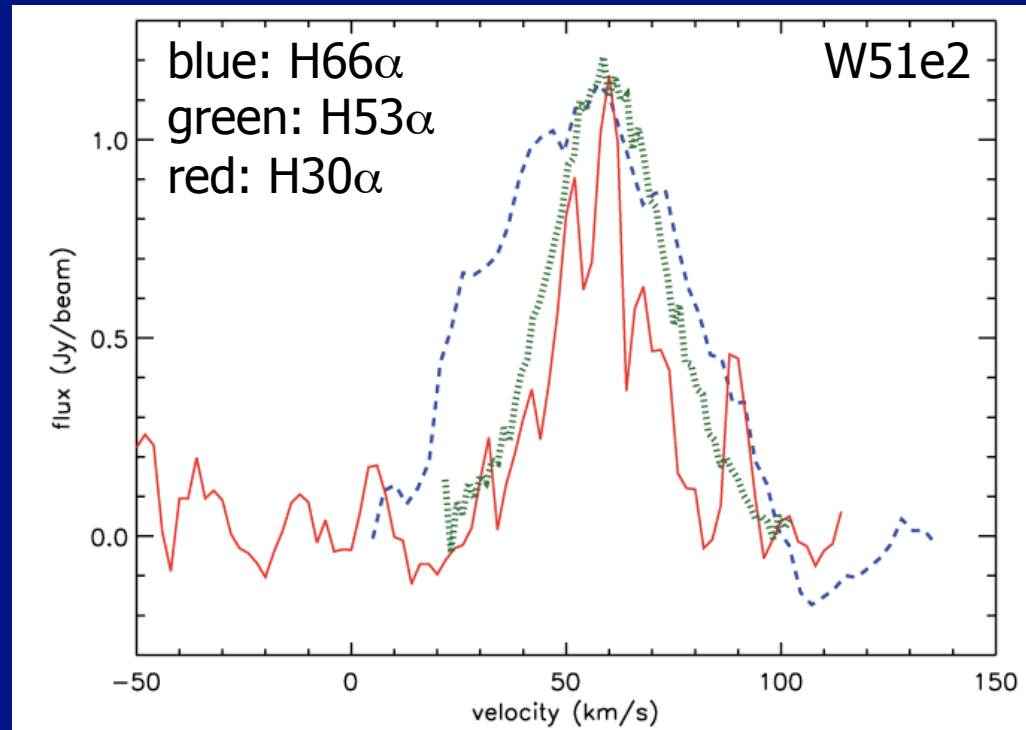


Recombination lines VII

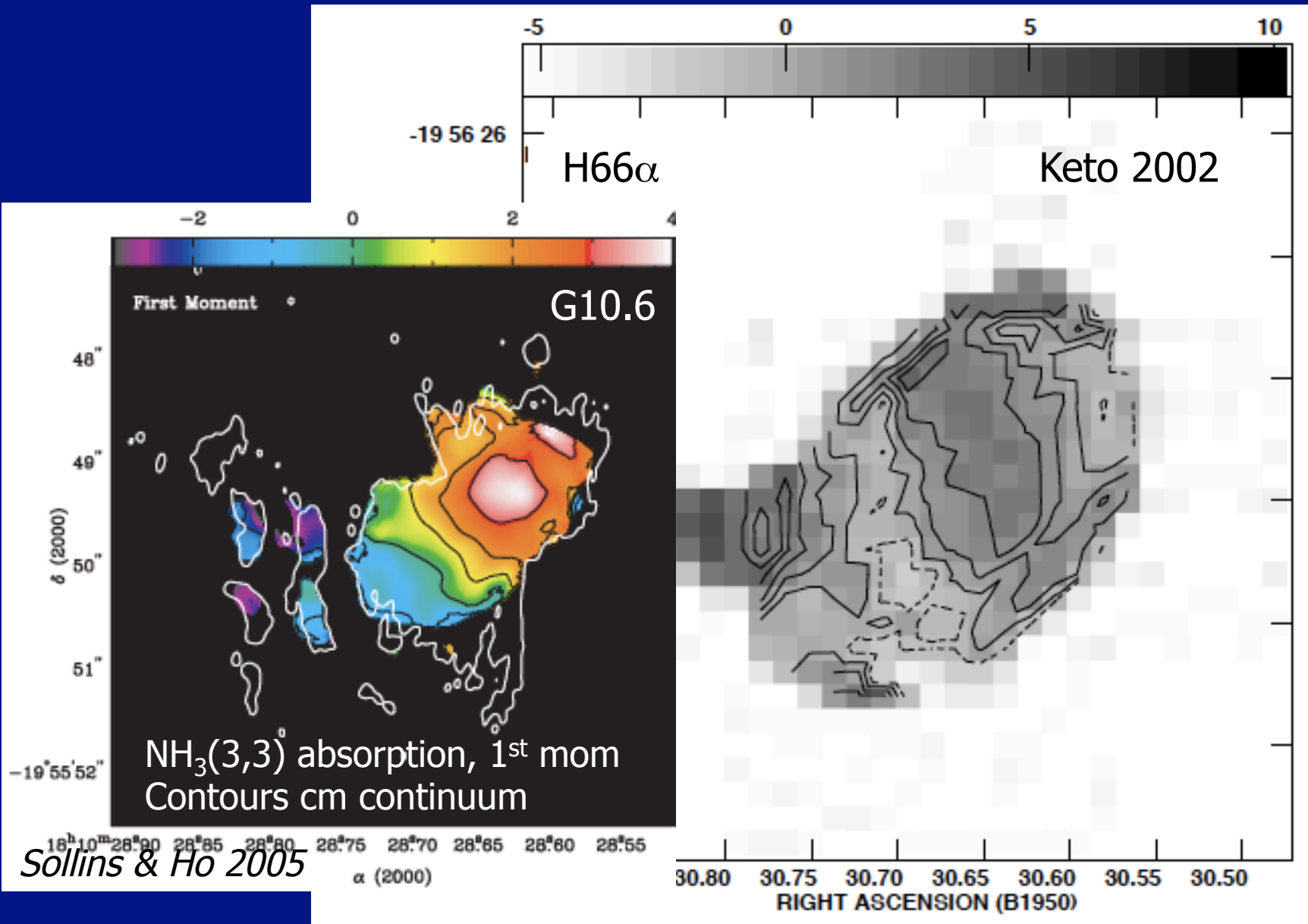
- Line strength, Einstein A:

$$A_{n+1,n} \approx \frac{64\pi^6 m_e e^{10}}{3 c^3 h^6 n^5} \approx 5.3 \times 10^9 \left(\frac{1}{n^5} \right) \text{ s}^{-1}$$

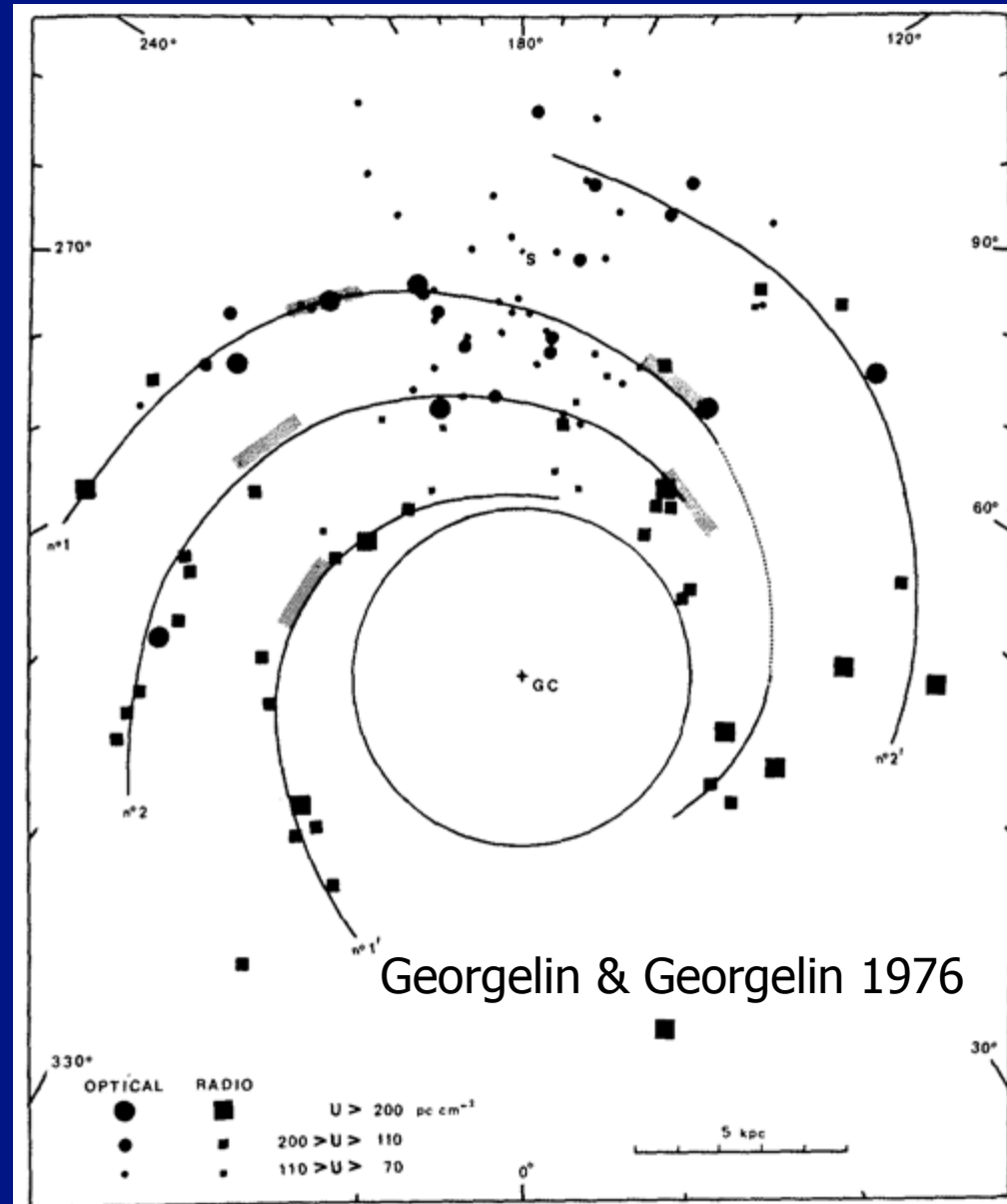
- Line width:
 - Natural: negligible
 - Thermal: Can be estimated from Maxwell distribution at given temperature:
For HII region with $T \sim 10000\text{K}$
 $\rightarrow \Delta v_{\text{therm}} \sim 20\text{km/s}$
 - Turbulent linewidth can be of similar magnitude
 - Collisional/pressure Stark broadening: Important for large n where interaction between atoms becomes significant.



Velocity pattern in HII regions



Velocity structure of the Milky Way



Georgelin & Georgelin 1976

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