

## Planet Formation

Nebula Hypothesis: Kant (1755) - Laplace (1796)

Planets formed from a flat nebula

(Laplace: Matter thrown off the Sun after it had condensed in its current state)

Nebular Hypothesis was revived by von Weizsäcker (1944) and Kuiper (1951): Planets in the Solar System formed from a slowly rotating cloud of dust and gas.

Quantitative Description for terrestrial planets:

Victor Safronov (1969): "Evolution of the Protoplanetary Cloud and Formation of the Earths and the Planets"

### Today

Gas-rich protoplanetary disks ( $\tau \sim 3 \text{ Myr}$ )

Gas-poor debris disks (late stages of planet formation)

### Observational Constraints:

#### a) Solar System

- Temporal and cosmochemical evolution
- Distribution of mass and angular momentum (concept of minimum mass solar nebula)
- Core masses of Jupiter & Saturn  
15-20  $M_{\oplus}$  for Saturn; 7-20  $M_{\oplus}$  for Jupiter
- Co-planar orbits; circular; same direction
- Distribution of minor bodies (asteroid belt, Kuiper belt)
- Small mass of moons relative to Earth & Venus

## g) Extrasolar Planets

Large diversity of planetary systems and exoplanet properties (eases of different detection techniques)

- Hot Jupiters & Neptunes (migration / in-situ)
- High frequency of Super-Earths / Mini-Neptunes
- Massive Jupiters around A-type stars
- Planets on eccentric orbits
- Planet pairs locked into dynamical mean-motion resonances (Early diff. migration via planet-disk interactions)
- Planets on retrograde and inclined orbits (dynamical interaction)
- Planets in binaries
- Strong correlation of stellar metallicity with gas giant formation

## Two scenarios

a) Core accretion - gas capture theory

(Sahrovov 1960, Mizuno 1980, Lissauer, Wetherill, Weidenschilling, Kley, ...)

b) Gravitational instability theory

(Kuiper 1951, Cameron 1962, Boss, Oerisen, ...)

## Gravitational instability theory

Remember Jeans criterion:

Density fluctuations in isothermal medium become unstable if its  $\lambda$  fulfills  $\lambda^2 > \pi c_s^2 / G \rho$

Sahrovov (1960) and Toomre (1964) refined the Jeans condition in a flat disk, including diff. rotation, gravity, and pressure effects.

$$Q = (c_s \kappa / \pi G \sigma) < 1 \text{ (instability)}$$

[local instability to axisymmetric perturbations]

## General idea for early calculations

Spherically symmetric condensations of approximately Jovian mass and solar composition formed.

→ Calculations involved the solution of the standard equations of stellar structure including radiative and convective energy transport and grain opacities: Contraction timescales depend on grain opacities ( $\sim 10^6$  yr)

If  $T > 2000K$  is reached  $\rightarrow$  hydrodynamic collapse on a timescale of 1 yr

- Boss (1998) indicated that a protoplanet of  $1 M_j$  could indeed form a core of a few  $\oplus$  in time with estimates of the core mass of Jupiter
- Only recently full numerical simulations of gravitational instability in disks with necessary spatial resolution  $\rightarrow$  mostly massive disks needed

Other issues

- a) Disk must not only be gravit. unstable, but it must be able to cool efficiently ( $\tau_{cool} < 3 \Omega_k^{-1}$  - Gammie 2001)
- b) Giant gaseous protoplanets would be formed in short timescales of about 1000 yrs
- c) Gas giants with  $\sim 10$  Jupiter masses at 100 AU could be formed; challenge is rapid mass transport inwards by spiral arms
- d) No natural explanation for metallicity - giant planet frequency relation because star determines radiative losses

## Core accretion - gas capture theory

Formation of a swarm of planetesimals and accretion on a planetary embryo in a certain feeding zone (Safronov 1969)

### Other assumptions

Total material budget in planet-forming zone was conserved and converted into the final planet mass:

Concept of "Minimum mass solar nebula"

(The solid-surface density just sufficient to correspond to the solid mass of final planets)

Growth time for Earth:  $100 \cdot 10^6$  yrs for Earth  
(consistent with estimates for growth time for Earth)

### Wehwill (1980):

$\sim 100$  lower-mass objects with lower spread over the terrestrial planet zone  $\rightarrow$  same timescale with approximately the correct number of objects

( $\rightarrow$  Mass problem: too low mass compared to Earth & Venus)



## Core accretion scenario (modern concept)

a) Dust particles grow to meter-sized bodies.

I:  $\mu\text{m}$ -sized grains grow through mutual collisions to fluffy fractal-type aggregates ( $m \propto D^{1.3}$ )

Relative velocities: Brownian motion, radial drift, sedimentation, turbulence)

II: cm-sized grains reach collis. velocities of  $10\text{ m/s}$  which leads to compaction and fragmentation

b) Meter-sized bodies decouple from gas and would migrate in  $100\text{ yrs}$

- Reduction of migration speed in turbulent flow with pressure bumps
- Gravitational instability in dead sublayers forms planetesimals ( $\sim 10\text{ km} - 100\text{ km}$ )

c) Runaway growth

Planetary embryos form by gravitational "clearing" of the planetesimal feeding zone.

Additional important factor in cross section

$$\pi b^2 = \pi r^2 (1 + v_{\text{esc}}^2 / \sigma^2)$$

$v_{\text{esc}}$  - escape velocity;  $\sigma$  - relative velocities

$$\frac{dm}{dt} = S_{\text{planet}} \cdot \sigma \pi b^2 \approx Z_{\text{planet}} \cdot \Omega \pi r^2 (1 + v_{\text{esc}}^2 / \sigma^2)$$

$$t_p = \sigma / \Omega$$

$$S_{\text{planet}} = Z_{\text{planet}} / t_p = \frac{Z_{\text{planet}} \cdot \Omega}{\sigma}$$

$$\Theta = v_{\text{esc}}^2 / \sigma^2 - \text{Salpeter factor}$$

At the beginning of the process  $\sigma$  small and  $\Theta \gg 1$

$$\frac{dm}{dt} \approx \bar{\Sigma}_{\text{planet}} \Omega \pi r^2 \Theta = \bar{\Sigma}_{\text{planet}} \cdot \pi \Omega r \frac{4GM}{\sigma^2}$$

$$\text{and } r \sim m^{1/3}$$

$$\boxed{\frac{dm}{dt} \sim m^{4/3}}$$

$m(t)$  diverges to infinity with time; formation of planetary embryos in feeding zone which dominate velocity dispersion: "oligarchs"

d) Self-regulated growth

Runaway ends when core (oligarch) dominates velocity dispersion and material in Hill

Sphere is depleted ( $r_H = a (m/3M_*)^{1/3}$ )

Maximum mass of a single body that has consumed all of the planetesimals in its Hill

Sphere  $\rightarrow$  Isolation mass

$$M_{\text{iso}} \sim \bar{\Sigma}_{\text{planet}}^{3/2} M_*^{-1/2} a^3$$

$$a = 1 \text{ AU} \quad \bar{\Sigma}_p = 10 \text{ g/cm}^2 \quad \rightarrow \sim 0.1 M_{\oplus}$$

$$a = 25 \text{ AU} \quad \rightarrow \quad 9 M_{\oplus}$$

Isolation mass played role for terrestrial planets

For giant planets may or may not have played a role depending on disk model

e) Grav. scattering and random collisions lead to the formation of terrestrial planets with timescales 10-100 Myr

### Giant planet formation

- Mizuno (1980) found that for a critical mass of about  $10 M_{\oplus}$  there is no hydrostatic equilibrium (Full calculation by Bodenheimer & Pollack 1986)
- Final mass determined by gas reservoir and how fast gas can be accreted

### Phases

- a) Core formation ( $\sim 1$  Myr)
- b) Hydrostatic growth (slow accretion of gas) ( $\sim 10$  Myr)
- c) Rapid growth when critical core mass is reached and  $\dot{M}_{\text{core}} \geq \dot{M}_{\text{gas}}$ ; rate of accretion accelerates dramatically
- d) Gas accretion reaches limiting value; equilibrium region of protoplanet contracts inside the effective accretion radius
- e) Accretion stops  $\rightarrow$  Planet enters isolation stage and contracts / cools at constant mass to its present stage



## Challenges

- a) One has to produce relatively massive core
- b) There must still be enough gas in the disk
- c) Migration into the central star

## Modifications

- a) Reduced opacities reduces formation timescale and one can form Jupiters with a lower mass core in the disk lifetime
- b) Reducing planetesimal accretion allows earlier runaway envelope

## Planet Migration

Type I: Acts on small planets which excite density waves at Lindblad resonances  
 waves interior to planet: positive torques  
 waves exterior to planet: negative torques

Sum of torques is negative  $\rightarrow$  inward migration

0.2 Myr for 1 Earth at 5 AU

(Same torques also ~~also~~ damp eccentricity)

No significant perturbations of density profile

Type II: More massive planets open a gap:

(linearity breaks down when  $\pi p_1 / M_* > (H/r)^3$   
 $\Rightarrow \sim 30 \text{ Earth}$ )

Gap structure depends on  $\pi p_1$ ,  $\alpha$ , disk structure

Planet is locked to disk evolution (viscous evol.)

Type III: Runaway migration associated with co-orbital torques

Lindblad resonances correspond to locations in the disk where the perturbing potential's frequency in the matter frame matches the epicyclic frequency

$$\omega(r) = m |\Omega_p - \Omega(r)| = \pm \kappa(r)$$

## Giant Impacts

- Very late in the formation of solar system, collisions between big protoplanets/embryos can take place leading to giant impacts

### a) Earth-Moon system

- Moon/Planet mass ratio bigger than for other satellite/planet configurations
- Low density of Moon  $\rightarrow$  Shattered Earth crust
- Impact of Moon-sized body required

### b) Rotational axis of Uranus

- Tilted over by impact?

### c) Chem. composition of Mercury

- High density, hardly any rock mantle
- Giant impact removed the mantle?

## Debris disks

- After  $\sim 10^7$  yr most gas-rich protoplanetary disks are gone  
(Removal processes)
- Dust grains would be removed from the system by radiative pressure and drag (Poynting-Robertson)
- Tiny, but measurable mass of solids detected ("Vega-like" stars)  
Example:  $\beta$  Pic  $\rightarrow$  carries also a planet

In this phase: Orbits planar and circular  
Accretion and ejection of remaining planetesimals;  
Belt structures in debris disks